

A new solution of Ernst's equation

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

1978 J. Phys. A: Math. Gen. 11 L75

(<http://iopscience.iop.org/0305-4470/11/4/002>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 129.252.86.83

The article was downloaded on 30/05/2010 at 18:48

Please note that [terms and conditions apply](#).

LETTER TO THE EDITOR

A new solution of Ernst's equation

C Hoenselaers

Raman Research Institute, Bangalore-560006, India

Received 15 February 1978

Abstract. A three-parameter solution of Ernst's equation, unfortunately not asymptotically flat, is presented. For special values of the parameters, the solution reduces to the Curzon solution.

In the elegant formulation of the problem of stationary gravitational fields by Ernst (1968, 1974) the field equations reduce to

$$\operatorname{Re}(\mathcal{E})\nabla^2\mathcal{E} = \nabla\mathcal{E} \cdot \nabla\mathcal{E}.$$

A survey of known solutions of this equation was given by Kinnersley (1975). Since then generalisations of the Tomimatsu-Sato (TS) solutions have been found by Yamazaki (1977a, b) and Cosgrove (1977). Asymptotically non-flat solutions corresponding to contractions of the TS metrics have been obtained by Ernst (1977).

The solution I have found is given in the usual spherical coordinates, r, θ as follows:

$$\mathcal{E} = \frac{A + iB}{N},$$

where

$$N = p^2 + q^2 r^4 \sin^4 \theta \exp(2m/r)$$

$$A = q^4 r^8 \sin^8 \theta \exp(3m/r) + p^2 r^2 \sin^2 \theta \exp(m/r) [2r^2 q^2 (1 - 8 \cos^2 \theta) + 8qr \cos \theta (2mq \cos \theta - a) + 4qm \cos \theta - 4q^2 m^2 \cos^2 \theta - a^2] + p^4 \exp(-m/r)$$

and

$$B = \exp(2m/r) \{ 16pq^3 \sin^2 \theta \cos^2 \theta r^4 + 8pq^2 r^3 \sin^2 \theta \cos \theta (a - 2q \cos \theta) + pqr^2 [-4q^2 m \cos^4 \theta + \cos^2 \theta (4qm + 6p^2 - a^2) - 4pqr \cos \theta \sin^2 \theta + a^2 - 2q^2] + 2q^2 r [a \cos \theta - pqm (3 \cos^2 \theta - 1)] + 2q^2 m \cos \theta (pqm - a) + 2p^3 qr^2 (3 \cos^2 \theta - 1) + 2p^2 r [a \cos \theta - pqm (3 \cos^2 \theta - 1)] + 2p^2 m \cos \theta (pqm - a) \}.$$

In this solution m and a are arbitrary real parameters, while p and q are related by $p^2 + q^2 = 1$.

For $q = 0$, $a = 0$, B vanishes and A reduces to the familiar expression for the Curzon metric.

As A/N and B/N behave like r^4 and r^0 , respectively, for large r , the solution cannot be asymptotically flat. The dominator N vanishes nowhere. One may thus suspect that the solution is free of singularities.

The author wishes to thank the Raman Research Institute for the award of a fellowship.

References

- Cosgrove C M 1977 *J. Phys. A: Math. Gen.* **10** 1481–524
Ernst F J 1968 *Phys. Rev.* **167** 1175–8
— 1974 *J. Math. Phys.* **15** 1409–12
— 1977 *J. Math. Phys.* **18** 233–4
Kinnorsley W 1975 *Proc. 7th Int. Conf. on General Relativity and Gravitation* eds G Shaviv and J Rosen pp 109–35
Yamazaki M 1977a *Prog. Theor. Phys.* **57** 1951–7
— 1977b *J. Math. Phys.* **18** 2502–8